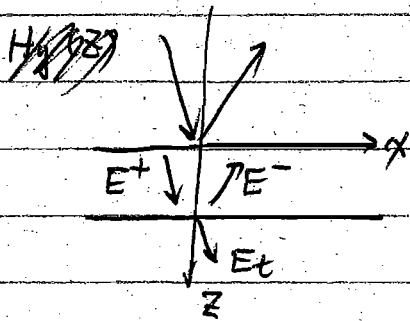


Review of Lecture 9



(1) One layer of film

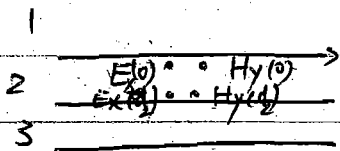
resultant wave method.

$$r = \frac{r_{12} + r_{23} e^{2i\phi_2}}{1 + r_{12} r_{23} e^{2i\phi_2}}$$

$$\phi_2 = \frac{\cos \theta_2 d N_2 \omega}{c_0} = \frac{2\pi \cos \theta_2 N_2}{\lambda_0}$$

Interference & tunneling phenomena.

(2) Multilayer



$$\begin{pmatrix} E_x(z) \\ H_y(z) \end{pmatrix} = \begin{pmatrix} \cos \phi_2 & i\beta_2 \sin \phi_2 \\ \frac{i}{\beta_2} \sin \phi_2 & \cos \phi_2 \end{pmatrix} \begin{pmatrix} E_x(0) \\ H_y(0) \end{pmatrix}$$

$$\begin{pmatrix} E_x(d) \\ H_y(d) \end{pmatrix} = \begin{pmatrix} \cos \phi_2 & -i\beta_2 \sin \phi_2 \\ -\frac{i}{\beta_2} \sin \phi_2 & \cos \phi_2 \end{pmatrix} \begin{pmatrix} E_x(0) \\ H_y(0) \end{pmatrix}$$

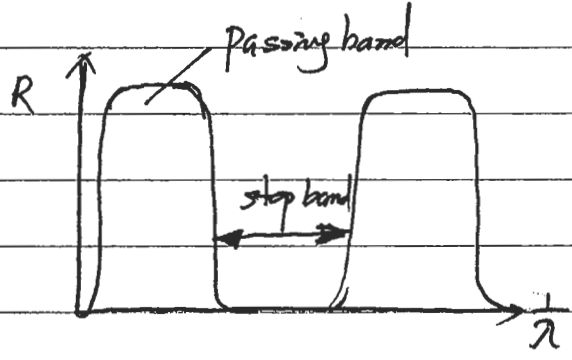
$$= M_2 \begin{pmatrix} E_x(d) \\ H_y(d) \end{pmatrix}$$

$$= M_2 \begin{pmatrix} E_x(z_3=0) \\ H_y(z_3=0) \end{pmatrix} = M_2 M_3 \begin{pmatrix} E_x(d_3) \\ H_y(d_3) \end{pmatrix}$$

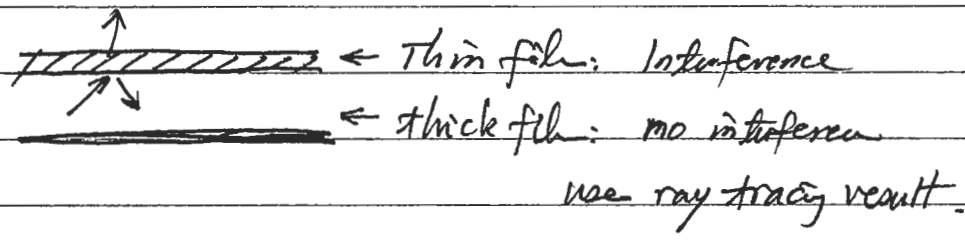
$$= M_2 \cdots M_n \begin{pmatrix} E_x(d_n) \\ H_y(d_n) \end{pmatrix}$$

$$A \begin{pmatrix} E_i \\ E_r \end{pmatrix} = M_2 \cdots M_n B E_t$$

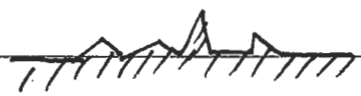
$$A = \begin{pmatrix} 1 & 1 \\ \frac{1}{\beta_1} & -\frac{1}{\beta_1} \end{pmatrix} \Rightarrow \text{Solve for } E_t/E_i, E_r/E_i$$



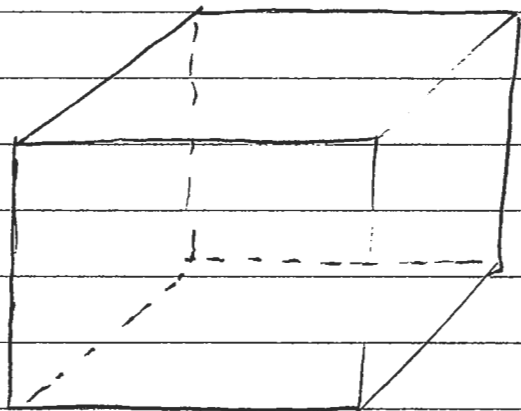
- * transition to thick films
- * combination of thin & thick films



* roughness interfaces.



* ~~Planck's law~~ Planck's law: a not so rigorous proof.



Consider a cube of length L
perfectly reflecting surface
large \Rightarrow free space

$$k_x^2 + k_y^2 + k_z^2 = \frac{\omega^2}{c^2}$$

$$k_x = \pm \frac{2\pi}{\lambda_x} = \pm \frac{2\pi n_x}{L}$$

$$n_x \lambda_x = L$$

Similarly k_y, k_z .

each combination of n_x, n_y, n_z is a different wave
at ω , there are many combinations

Quantum: each wave must have an energy of
 $E_n = n \hbar \omega$ ($n=0, 1, \dots$)

Statistical physics: $A \exp\left(-\frac{E_n}{kT}\right)$ = probability of actually
having E_n .
 \uparrow Boltzmann constant

$$\sum_n A \exp\left(-\frac{n \hbar \omega}{kT}\right) = 1$$

$$A = \frac{1}{1 - \exp\left(-\frac{\hbar \omega}{kT}\right)}$$

Probability of each wave \Rightarrow average # of photons a wave has

$$\sum n \cdot p(n) = \frac{1}{e^{\frac{\hbar \omega}{kT}} - 1} = f(\omega)$$

The average energy of each wave

$$\hbar \omega f(\omega)$$

We need to find out how many waves we have at ω for a length of L .

Volume per mode n_y, n_x, n_z change by ± 1 . $(\frac{2L}{\lambda})^3$

$$\# \text{ of mode} = \frac{4\pi dk_x dk_y dk_z}{(2\pi/L)^3} \cdot 2$$

$$\begin{aligned} \text{density of states } D(\omega) &= 2 \cdot \frac{1}{d\omega} \times \frac{L^3}{8\pi^3} \cdot dk_x dk_y dk_z \\ &= 2 \frac{V}{8\pi^3} \frac{4\pi k^2 dk}{d\omega} \end{aligned}$$

$$k = \frac{\omega}{c} \quad \omega = ck$$

$$D(\omega) = \frac{V \omega^2}{c^3 \pi^2}$$

so, energy density per unit volume

$$U(\omega) = \hbar \omega f(\omega) D(\omega) = \frac{\hbar \omega^3}{\pi^2 c^3 (e^{\frac{\hbar \omega}{kT}} - 1)}$$

$$\text{Intensity } \frac{c}{4\pi} U(\omega) = I(\omega) = \frac{\hbar \omega^3}{4\pi^3 c^2 (e^{\frac{\hbar \omega}{kT}} - 1)}$$